

Laplace Transform: Developing the Variational Iteration Method

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Abstract:

The classification of the Lagrange multiplier plays an important role in the variational iteration method and the variational theory is widely used for this purpose. This paper suggests an easier approach by the Laplace transform to determine the multiplier, making the process obtainable to researchers facing different nonlinear problems. A nonlinear oscillator is adopted as an illustration to elucidate the detection process and the solution process, only one iteration leads to an ultimate result.

Introduction

The variational iteration method was proposed in late 1990s to solve a scalar flow with fractional derivatives and a nonlinear oscillator [1, 2], and this method has widely used as a main mathematical tool to solving various nonlinear equations. Due to general study of the method by numerous authors, for examples, Ji-Huan He [3–5], D.D. Ganji [6], T. Ozis and A. Yildirim [7], M.A. Noor and S.T. Mohyud-Din [8], it has completely developed into a fully fledged method in mathematics. Using “variational iteration method” as a searching topic in Clarivate’s web of science, we found 3761 hits on 24 November 2018. The identification of the Lagrange multiplier in the method requires the facts of the variational theory [8–11], and the complex detection process might delay applications of the method to practical problems. This paper suggests an easier detection process by the Laplace transform, which is available in all mathematics textbooks.

The identification of the Lagrange multiplier by the Laplace transform

Consider a general non-linear oscillator equation in the form:

$$u''(t) + f(u) = 0 \dots\dots\dots (1)$$

with initial conditions $u(0) = A, \quad u'(0) = 0$
(2)

We can rewrite Eq. (1) as

$$u'' + \omega^2 u + g(u) = 0 \dots\dots\dots (3)$$

where ω is the frequency to be auxiliary determined, $g(u) = f(u) - \omega^2 u$.

According to the variational iteration method (VIM), the alteration functional for Eq. (3) is given as [1–5]

$$u_{n+1}(t) = u_n(t) + \int_0^t \lambda(t, \xi) [u_n''(\xi) + \omega^2 u_n(\xi) + \tilde{g}(u_n)] d\xi, \quad n = 0, 1, 2, \dots$$

.....(4)

where λ is a general Lagrange multiplier, and it can be optimally resolved from the stationary conditions of Eq. (4) with respect to u_n using the variational theory [9–11]. The subscript n represents the n th approximation and \tilde{g} is a constrained variation, i.e., $\delta \tilde{g} = 0$. There are many publications discussing how to recognize the multiplier effectively.

Hereby we will show a different approach to the classification of the multiplier. Starting from some pioneering thoughts going back to Abassy, El-Tawil and El-Zoheiry in 2007 [12], Mokhtari and Mohammadi in 2009 [13], Hesameddini and Latifizadeh in 2009 [14], the Laplace transform was adopted in the variational iteration method. Abassy, El-Tawil and El-Zoheiry [12] used Laplace transform in the result method, the variational iteration technique leads to a succession of linear equations, which can be easily solved by the Laplace transform. Mokhtari and Mohammadi [13] found with the intention of the variational iteration algorithm could be simply constructed by the Laplace transform without using the alteration functional (the variational theory) and restricted variations. Hesameddini and Latifizadeh [14] found that Laplace transform could erect iteration algorithms as those by the variational iteration method. When solving a fractional differential equation, the variational iteration method shows some obvious advantages over others [15–19], and the Laplace transform plays an even more important role in the solution process [20–22]. The present method of below is also applicable for fractal derivative equations [23–27].

Generally the Lagrange multiplier can be expressed in the form [1–5]

$$\lambda = \lambda(t - \xi) \dots\dots\dots (5)$$

In view of Eq. (5), the correction functional given in Eq. (4) is basically the convolution; hence we can use the Laplace transform easily.

Applying the Laplace transform on both sides of Eq. (4), the correction functional will be transformed in the following manner

$$\begin{aligned} L[u_{n+1}(t)] &= L[u_n(t)] + L\left[\int_0^t \lambda(t - \xi) [u_n''(\xi) + \omega^2 u_n(\xi) + \tilde{g}(u_n)] d\xi\right] \\ &= L[u_n(t)] + L[\lambda(t) * (u_n''(t) + \omega^2 u_n(t) + \tilde{g}(u_n))] \\ &= L[u_n(t)] + L[\lambda(t)] L[u_n''(t) + \omega^2 u_n(t) + \tilde{g}(u_n)] \\ &= L[u_n(t)] + L[\lambda(t)] [(s^2 + \omega^2)L[u_n(t)] - s u_n(0) - u_n'(0) + L[\tilde{g}(u_n)]] \end{aligned}$$

$$\dots\dots\dots (6)$$

The optimal value of λ can be obtained by making Eq. (6) stationary with respect to $u_n(t)$, this requires

$$\begin{aligned} \frac{\delta}{\delta u_n} L[u_{n+1}(t)] &= \frac{\delta}{\delta u_n} L[u_n(t)] + \frac{\delta}{\delta u_n} L[\lambda(t)] [(s^2 + \omega^2)L[u_n(t)] - su_n(0) - u'_n(0) + L[\tilde{g}(u_n)]] \\ &= \{1 + L[\lambda(t)](s^2 + \omega^2)\} \frac{\delta L[u_n(t)]}{\delta u_n} = 0 \end{aligned} \dots\dots\dots (7)$$

From Eq. (7), we have

$$L[\lambda] = -\frac{1}{(s^2 + \omega^2)} \dots\dots\dots(8)$$

In the above derivation, we assume that

$$\frac{\delta L[\tilde{g}(u_n)]}{\delta u_n} = 0 \dots\dots\dots(9)$$

The inverse Laplace transform for Eq. (8) results in

$$\lambda(t) = -\frac{1}{\omega} \sin \omega t \dots\dots\dots(10)$$

We, therefore, identify the Lagrange multiplier much easier than that by the variational theory.

An example

As an example, we consider the following oscillator [28,29]:

$$(1 + \alpha u^2)u'' + \alpha uu'^2 - u(1 - u^2) = 0 \dots\dots\dots (11)$$

with initial conditions

$$u(0) = A, \quad u'(0) = 0 \dots\dots\dots (12)$$

This example was solved by the homotopy perturbation method [28]. We write Eq. (11) in the form

$$u'' + \omega^2 u + g(u) = 0 \dots\dots\dots (13)$$

Where

$$g(u) = -(1 + \omega^2)u + \alpha u^2 u'' + \alpha u u'^2 + u^3$$

We have the following iteration formula

$$\begin{aligned} L[u_{n+1}(t)] &= L[u_n(t)] - L\left[\int_0^t \frac{1}{\omega} \sin \omega(t - \xi) (u_n''(\xi) + \omega^2 u_n(\xi) + g(u_n)) d\xi\right] \\ &= L[u_n(t)] - \frac{1}{\omega} L[\sin \omega t] L[u_n''(t) + \omega^2 u_n(t) + g(u_n)] \\ &= L[u_n] - \frac{1}{\omega} L[\sin \omega t] L[u_n'' - u_n + \alpha u_n^2 u_n'' + \alpha u_n u_n'^2 + u_n^3] \\ &\dots\dots\dots (14) \end{aligned}$$

Assuming the initial solution is

$$u_0(t) = A \cos \omega t \dots\dots\dots (15)$$

we have

$$\begin{aligned} L[u_1(t)] &= L[A \cos \omega t] - \frac{1}{\omega} L[\sin \omega t] L\left[-A\omega^2 \cos \omega t - A \cos \omega t - \frac{\alpha A^3 \omega^2}{4} (3 \cos \omega t + \cos 3\omega t) \right. \\ &\quad \left. + \frac{\alpha A^3 \omega^2}{4} (\cos \omega t - \cos 3\omega t) + \frac{A^3}{4} (3 \cos \omega t + \cos 3\omega t)\right] \\ &= L[A \cos \omega t] - \frac{1}{\omega} \left(-A\omega^2 - A + \frac{3}{4}A^3 - \frac{1}{2}\alpha A^3 \omega^2\right) L[\sin \omega t] L[\cos \omega t] \\ &\quad - \frac{1}{\omega} \left(\frac{1}{4}A^3 - \frac{1}{2}\alpha A^3 \omega^2\right) L[\sin \omega t] L[\cos 3\omega t] \\ &\dots\dots\dots (16) \end{aligned}$$

The inverse Laplace transform on Eq. (16) results in the first order approximate solution:

$$\begin{aligned} u_1(t) &= A \cos \omega t - \frac{1}{\omega} \left(-A\omega^2 - A + \frac{3}{4}A^3 - \frac{1}{2}\alpha A^3 \omega^2\right) \left(\frac{1}{2}t \sin \omega t\right) \\ &\quad - \frac{1}{\omega} \left(\frac{1}{4}A^3 - \frac{1}{2}\alpha A^3 \omega^2\right) \left(\frac{1}{8\omega} (\cos \omega t - \cos 3\omega t)\right) \\ &\dots\dots\dots \\ (17) \end{aligned}$$

No secular-term in Eq. (17) requires that

$$\frac{1}{\omega} \left(-A\omega^2 - A + \frac{3}{4}A^3 - \frac{1}{2}\alpha A^3 \omega^2\right) = 0 \dots\dots\dots (18)$$

which leads to the following result

$$\omega = \sqrt{\frac{\frac{3}{4}A^2 - 1}{1 + \frac{1}{2}\alpha A^2}} \dots\dots\dots (19)$$

Eq. (19) is exactly the same as that obtained by the homotopy perturbation method [28] or He's frequency–amplitude formulation [30].

Discussion and conclusion

In this short paper we apply the Laplace transform to identify easily the Lagrange multiplier. As the Laplace transform is widely known to almost all non-mathematicians, such identification of the Lagrange multiplier makes the variational iteration method accessible to all researchers who face various nonlinear problems. The use of the variational iteration process now requires no particular knowledge of elusive calculus of variations. Though this paper gives a basic solution process to a nonlinear oscillator, the method is valid for other nonlinear problems as well.

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